

# Model-independent analysis of charged-lepton-flavour-violating $\tau$ processes

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## 1. Motivation

Neutrinos oscillate,  $\nu_\ell \leftrightarrow \nu_{\ell'}$ : why shouldn't charged leptons?

- Minimal extensions of the SM predict negligible flavour violation in the charged lepton sector (CLFV)
- New physics scenarios allow for enhanced CLFV

Already limits on the first and second family:

$$\mu N \rightarrow e N' \rightarrow R_{\mu e}^{Au} < 7 \times 10^{-13} \quad (\text{Sindrum II, 2006})$$

Belle and BaBar set limits on the third family for hadronic decays:

$$\tau \rightarrow \ell H(H) \quad (\text{e.g., } \Gamma_{\tau\mu\eta} < 6.5 \times 10^{-8})$$

Prospects:

Experiment **NA64** at CERN will look for  $\ell N \rightarrow \tau N'$ . They claim a sensitivity of:  $R_{\ell\tau} = \frac{\sigma(\ell + N \rightarrow \tau + X)}{\sigma(\ell + N \rightarrow \ell + X)} \sim 10^{-12} - 10^{-13}$ ,  $\ell = e, \mu$

**Belle II** will improve the limits for hadronic  $\tau$  decays  $\tau \rightarrow \ell H(H)$  at least **one order** of magnitude

## 2. Our project

Use of the **SMEFT** up to **D-6 operators** to analyse the  $\tau$ -involved processes

- $l - \tau$  conversion in nuclei:

$$l \mathcal{N}(A, Z) \longrightarrow \tau X$$

- Hadronic  $\tau$  decays:

$$\tau \rightarrow lP$$

$$\tau \rightarrow lPP \quad (l \neq \nu_\tau)$$

$$\tau \rightarrow lV$$

**Numerical analysis** with the experimental and theoretically expected limits on those decays by **NA64**, **Belle** and **Belle II**

# Dimension-6 Operators in the SMEFT

$$\mathcal{L}_{SM} = \mathcal{L}_{SM}^{(4)} + \frac{1}{\Lambda} \sum_k C_k^{(5)} \mathcal{Q}_k^{(5)} + \frac{1}{\Lambda^2} \sum_k C_k^{(6)} \mathcal{Q}_k^{(6)} + \mathcal{O}\left(\frac{1}{\Lambda^3}\right)$$

[Grzadkowski et al., 2010]

$$\mathcal{Q}_{Lu} = (\bar{L}_p \gamma_\mu L_r)(\bar{u}_s \gamma^\mu u_t) \xrightarrow{CLFV} (\bar{\tau} \gamma_\mu P_L \ell)(\bar{u}_s \gamma^\mu P_R u_t)$$

$$\mathcal{Q}_{\varphi e} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}_p \gamma^\mu e_r) \xrightarrow{CLFV} \frac{g_V^2}{2 \cos \theta_W} Z_\mu (\bar{\tau} \gamma^\mu P_R \ell)$$

General cross section of the process:

$$\sigma(\ell \mathcal{N} \rightarrow \tau X) = Z \sigma(\ell p \rightarrow \tau X) + (A - Z) \sigma(\ell n \rightarrow \tau X)$$

In terms of **partonic** cross sections at leading order in  $\alpha_s$ :

$$\begin{aligned} \sigma(\ell \mathcal{N}(P) \rightarrow \tau X) = & \sum_{i,j} \int_0^1 d\xi \{ \sigma[\ell q_i(\xi P) \rightarrow \tau q_j] f_i(\xi) \\ & + \sigma[\ell \bar{q}_j(\xi P) \rightarrow \tau \bar{q}_i] f_j(\xi) \} \end{aligned}$$

$f_i(\xi)$  are the nuclear parton distribution functions

[Kovařík et al., 2016]

$$\sigma[\ell q_i(\xi P) \rightarrow \tau q_j] = \sigma(C_k^{(6)}, \xi)$$

# Hadronic $\tau$ decays

We consider three different flavour violating **hadronic**  $\tau$  decays:

$$\textcircled{1} \quad \tau \rightarrow \ell P : \quad P = \pi^0, K^0, \bar{K}^0, \eta, \eta'$$

$$\textcircled{2} \quad \tau \rightarrow \ell PP : \quad PP = \pi^+\pi^-, K^0\bar{K}^0, K^+K^-, \pi^+K^-, K^+\pi^-$$

$$\textcircled{3} \quad \tau \rightarrow \ell V : \quad V = \rho^0, \phi, \omega, K^{*0}, \bar{K}^{*0}$$

The **quark currents** in the D-6 operators are hadronized through **Chiral perturbation theory** ( $\chi PT$ ) and **Resonance chiral theory** ( $R\chi T$ )

[Weinberg, 1979]

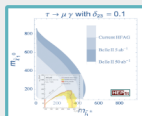
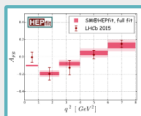
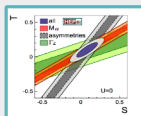
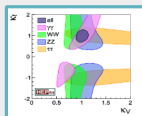
[Ecker et al., 1989]

$$Q_{Lu} \rightarrow (\bar{u}_s \gamma^\mu u_t) \longrightarrow \begin{cases} P & \chi PT + R\chi T \\ PP & \chi PT + R\chi T \\ V & R\chi T \end{cases}$$



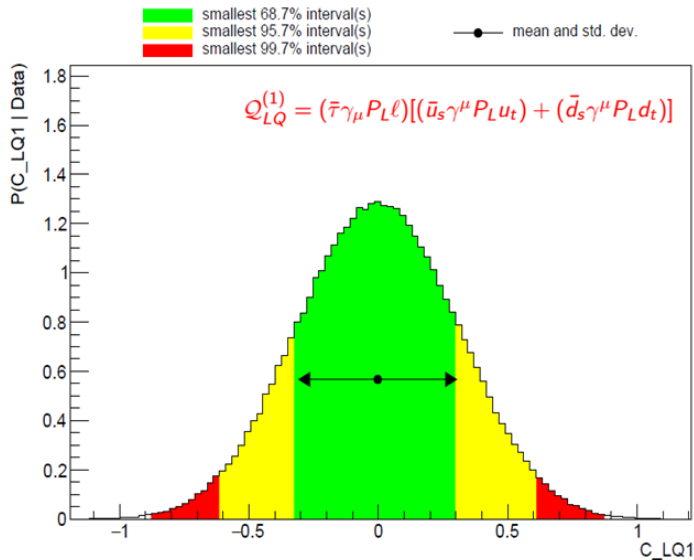
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**HEPfit**: a Code for the Combination of Indirect and Direct Constraints on High Energy Physics Models.



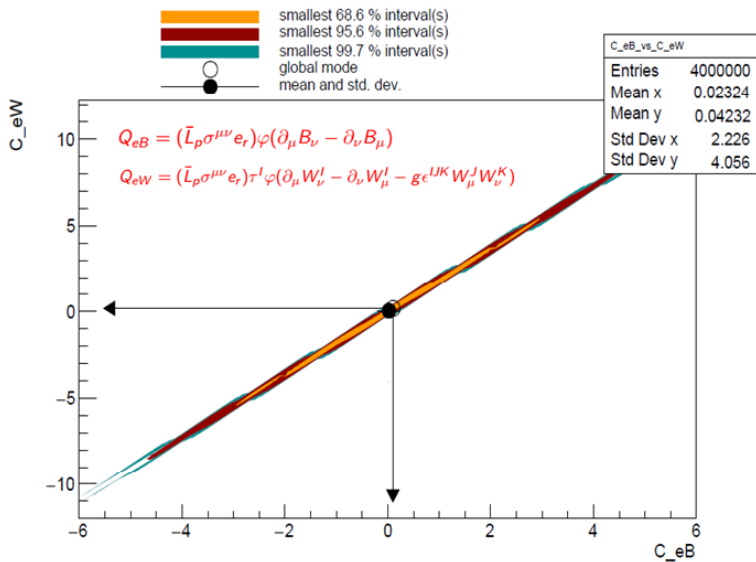
<http://hepfit.roma1.infn.it/>





$$\Lambda = 1 \text{ TeV}$$

$$\begin{aligned}
 C_{LQ}^{(1)} &= -0.002334 \pm 0.3044 & C_{LQ}^{(3)} &= 0.001823 \pm 0.5034 \\
 C_{eu} &= 0.002215 \pm 0.519 & C_{ed} &= -0.001052 \pm 0.3232 \\
 C_{Lu} &= 0.0129 \pm 0.828 & C_{Ld} &= 0.004037 \pm 0.5295 \\
 C_{Qe} &= 0.001388 \pm 0.4445 & C_{LedQ} &= 0.0153 \pm 0.7735 \\
 C_{LeQu}^{(1)} &= 0.007953 \pm 1.662 & C_{LeQu}^{(3)} &= 0.001444 \pm 0.1546 \\
 C_{\varphi L}^{(1)} &= -0.1673 \pm 11.53 & C_{\varphi e} &= 0.003808 \pm 0.6962 \\
 C_{\varphi L}^{(3)} &= 0.1646 \pm 11.5 & C_{eB} &= 0.02324 \pm 2.226 \\
 C_{eW} &= 0.04232 \pm 4.056 & &
 \end{aligned}$$



# Conclusions

Future experiments will set (if not observe) stronger limits in CLFV  $\tau$ -involved processes:

- NA64  $\longrightarrow \ell - \tau$  conversion in nuclei
- Belle II  $\longrightarrow$  Hadronic  $\tau$  decays

We performed a **Model-independent analysis** through the SMEFT up to D-6 operators of the whole basis of CLFV operators, then we set **constraints on their Wilson coefficients**





Some preliminary results have been shown

*Thank you very much for your attention!*



<http://lhcpheo.ific.uv-csic.es/>

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