

# **Review on hadronic B decays**

**Jike Wang (Wuhan University)**

on behalf of the LHCb, including results  
from ATLAS and CMS

# Outline of the talk

- New resonances:  $(\Lambda_b^0 \pi^+ \pi^-)$ ,  $(\Lambda_b^0 \pi^\pm)$
- Observation of  $B_{(s)}^0 \rightarrow J/\psi p \bar{p}$ ,  $\Lambda_b^0 \rightarrow \chi_{c1}(3872) p K^-$ ,  
 $B_s^0 \rightarrow \bar{D}^{*0} \phi$ , excited  $B_c^+$  state,  $B^0 \rightarrow (Z_c \rightarrow \eta_c(1S)) K^+ \pi^-$ ,  
 $\Lambda_b \rightarrow J/\psi \Lambda \phi$  (CMS)
- Amplitude analysis of  $B^\pm \rightarrow \pi^\pm K^+ K^-$ ,  $B_s^0 \rightarrow K^{*0} \bar{K}^{*0}$ ,  
 $B_s^0 \rightarrow K_s^0 K^\pm \pi^\mp$
- Several sources of CPV in  $B^+ \rightarrow \pi^+ \pi^+ \pi^-$ ; model independent  
observation of  $B^0 \rightarrow J/\psi K^+ \pi^-$ ; CPV phase  $\phi_s$  in  
 $B_s^0 \rightarrow J/\psi \phi$  (ATLAS)

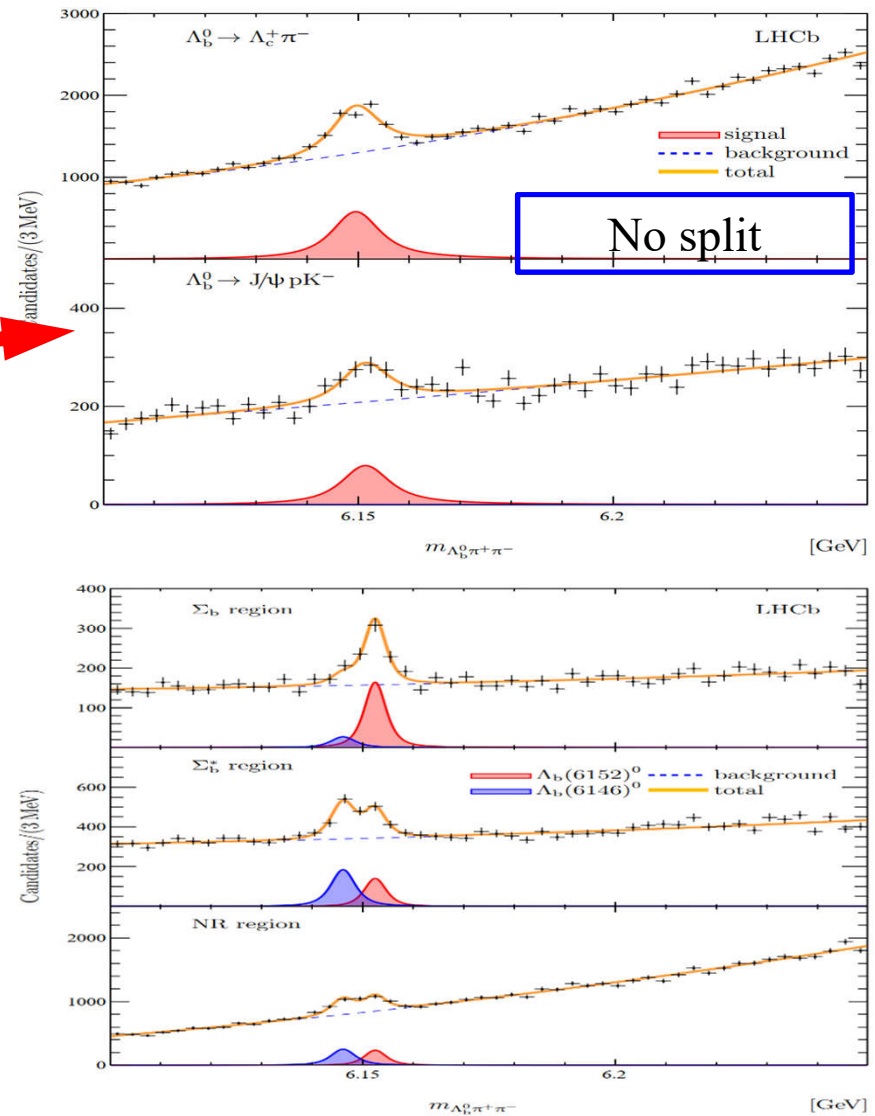
# New resonances in the $\Lambda_b^0 \pi^+ \pi^-$ system

- Excited states to the  $\Lambda_b^0 \pi^+ \pi^-$  has already been studied by the LHCb: e.g.  $\Lambda_b(5912)^0$
- More states are predicted in the mass region near/above 6.1 GeV ( $> \Sigma_b^{(*)\pm} \pi^\mp$ ).

## Split $m_{\Lambda_b^0 \pi^\pm}$ into 3 regions:

- Within the natural width of the known  $\Sigma_b^\pm$  mass
- Within the natural width of the known  $\Sigma_b^{(*)\pm}$  mass
- Remaining nonresonant (NR) region.

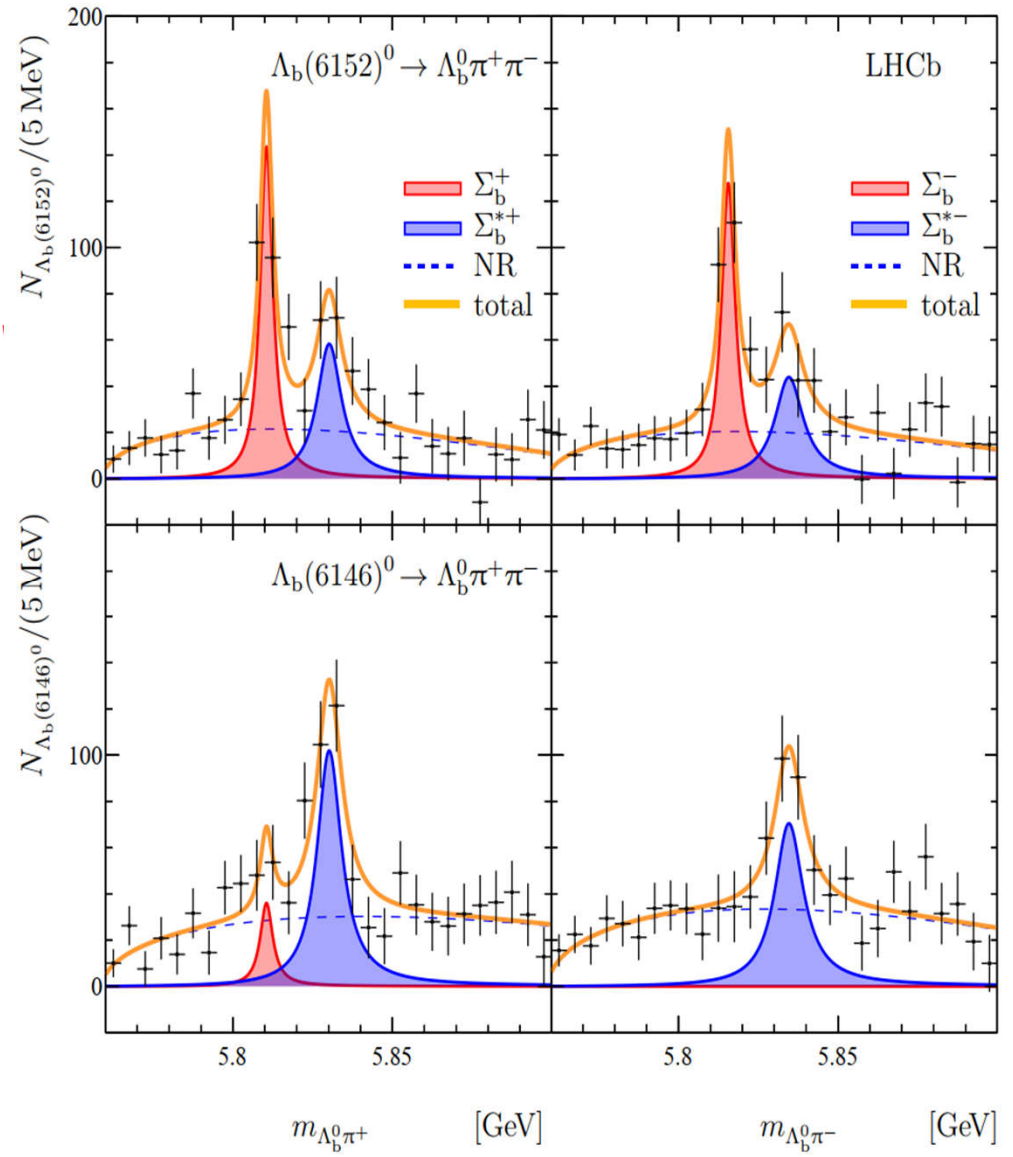
[ARXIV:1907.13598](https://arxiv.org/abs/1907.13598);  $9fb^{-1}$



**Look different  $\Rightarrow$  two narrow peaks**

# New resonances in the $\Lambda_b^0 \pi^+ \pi^-$ system

- Significant  $\Lambda_b(6152)^0 \rightarrow \Sigma_b^\pm \pi^\mp$  and  $\Lambda_b(6152)^0 \rightarrow \Sigma_b^{*\pm} \pi^\mp$  signals are observed
  - 1/3 and 1/4 of the signal decays in the sample
- $\Lambda_b(6146)^0 \rightarrow \Sigma_b^{*\pm} \pi^\mp$  decays account for about half of the observed decay rate
- Consistent with the predictions for the doublet of  $\Lambda_b(1D)^0$ 
  - Similar natural widths are expected for the two states of the doublet.
- Interpretation of these states as excited  $\Sigma_b^0$  states cannot be excluded.



[ARXIV:1907.13598 \(to PRL\)](https://arxiv.org/abs/1907.13598);  $9fb^{-1}$

# Two new resonances in the $(\Lambda_b^0 \pi^\pm)$ system

- Beyond these ground states  $(\Sigma_b^\pm, \Sigma_b^{*\pm})$ 
  - Radially/orbitally excited states are expected at higher masses, but only a few observed
- These excited states will cast light on the internal mechanisms governing the dynamics of the constituent quarks

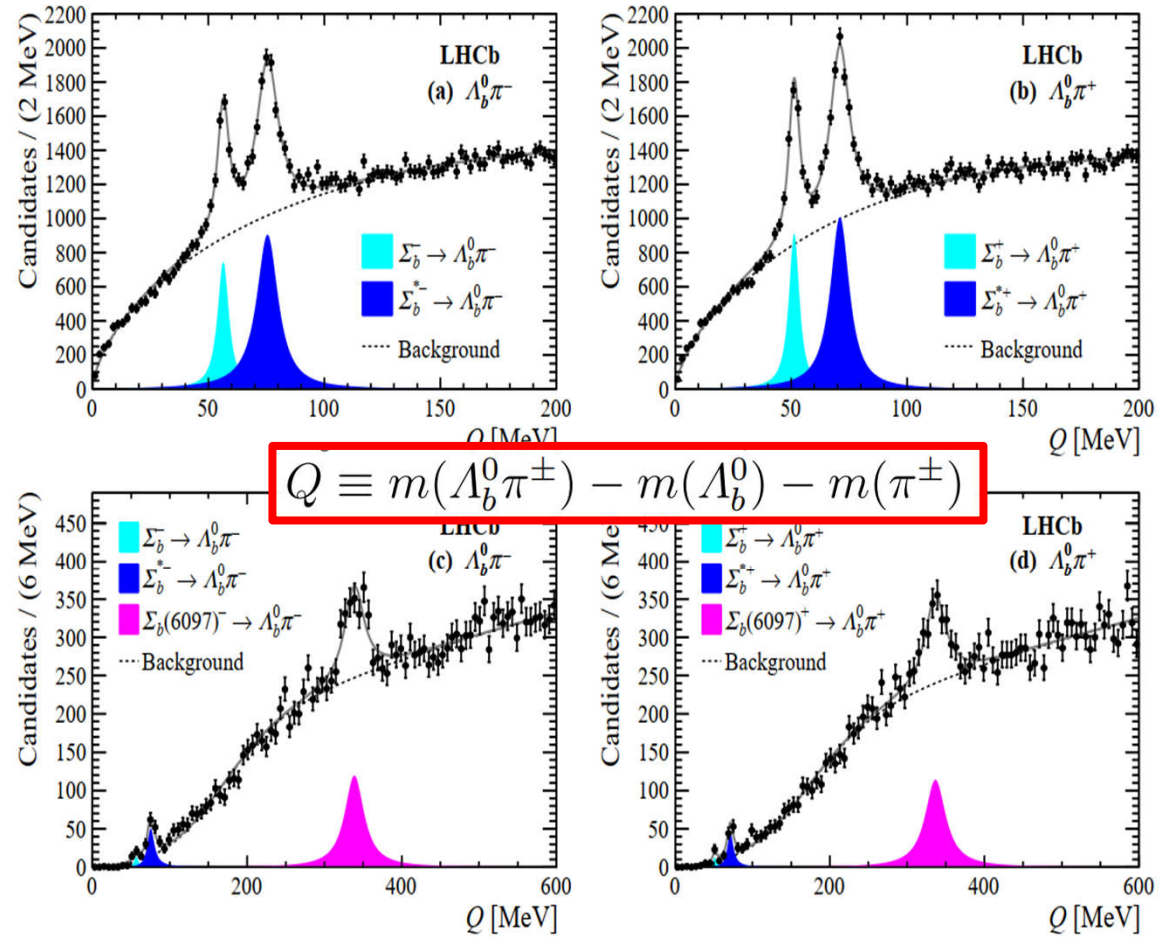


Figure 2: Mass distribution for selected  $\Lambda_b^0 \pi^\pm$  candidates. The points show experimental data. The left (right) column shows  $\Lambda_b^0 \pi^-$  ( $\Lambda_b^0 \pi^+$ ) combinations. The top row shows the fits to the lower-mass states  $\Sigma_b^\pm$  and  $\Sigma_b^{*\pm}$ . The lower row presents the fits to the new mass peaks with the requirement  $p_T(\pi_s^\pm) > 1000$  MeV.

**PRL. 122(2019) 012001; 9fb<sup>-1</sup>**

# The measurements of the $(\Lambda_b^0 \pi^\pm)$ resonances

- Observation of two new mass peaks in the  $\Lambda_b^0 \pi^+$  and  $\Lambda_b^0 \pi^-$  systems
- The ground-states ( $\Sigma_b^\pm$ ,  $\Sigma_b^{*\pm}$ ) are also confirmed and the masses and widths precisely measured.

Quantity	Value [MeV]
$m(\Sigma_b(6097)^-)$	$6098.0 \pm 1.7 \pm 0.5$
$m(\Sigma_b(6097)^+)$	$6095.8 \pm 1.7 \pm 0.4$
$\Gamma(\Sigma_b(6097)^-)$	$28.9 \pm 4.2 \pm 0.9$
$\Gamma(\Sigma_b(6097)^+)$	$31.0 \pm 5.5 \pm 0.7$
$m(\Sigma_b^-)$	$5815.64 \pm 0.14 \pm 0.24$
$m(\Sigma_b^{*-})$	$5834.73 \pm 0.17 \pm 0.25$
$m(\Sigma_b^+)$	$5810.55 \pm 0.11 \pm 0.23$
$m(\Sigma_b^{*+})$	$5830.28 \pm 0.14 \pm 0.24$
$\Gamma(\Sigma_b^-)$	$5.33 \pm 0.42 \pm 0.37$
$\Gamma(\Sigma_b^{*-})$	$10.68 \pm 0.60 \pm 0.33$
$\Gamma(\Sigma_b^+)$	$4.83 \pm 0.31 \pm 0.37$
$\Gamma(\Sigma_b^{*+})$	$9.34 \pm 0.47 \pm 0.26$
$m(\Sigma_b^{*-}) - m(\Sigma_b^-)$	$19.09 \pm 0.22 \pm 0.02$
$m(\Sigma_b^{*+}) - m(\Sigma_b^+)$	$19.73 \pm 0.18 \pm 0.01$
$\Delta(\Sigma_b(6097)^\pm)$	$-2.2 \pm 2.4 \pm 0.3$
$\Delta(\Sigma_b^\pm)$	$-5.09 \pm 0.18 \pm 0.01$
$\Delta(\Sigma_b^{*\pm})$	$-4.45 \pm 0.22 \pm 0.01$

- Compatible with being  $\Sigma_b(1P)$  excitations
- Other interpretations, such as molecular states, may also be possible

# Observation of $B_{(s)}^0 \rightarrow J/\psi p \bar{p}$ decays

- Sensitive to pentaquark in the  $J/\psi p$  and  $J/\psi \bar{p}$  components and to glueball states ( $p \bar{p}$ )
- Baryonic  $B_{(s)}^0$  decays:
  - Study the dynamics of the final baryon-antibaryon system and the threshold enhancement
- Theoretical expectation for the  $Br(B_{(s)}^0 \rightarrow J/\psi p \bar{p})$  is at the level of  $10^{-9}$ .
  - Intermediate pentaquark or glueball state can enhance

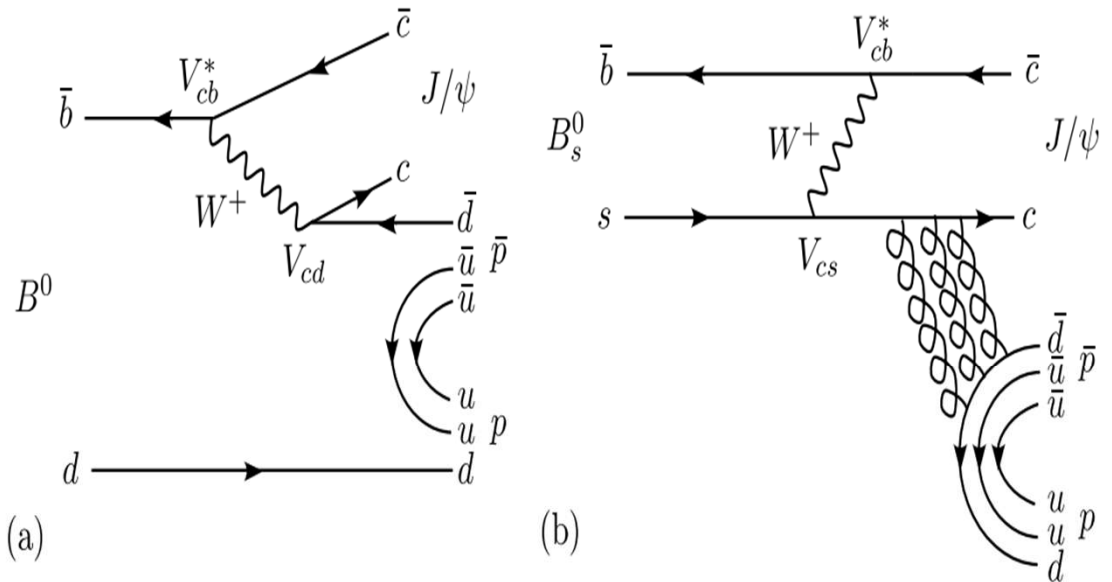


Figure 1: Leading Feynman diagrams for (a)  $B^0 \rightarrow J/\psi p \bar{p}$  and (b)  $B_s^0 \rightarrow J/\psi p \bar{p}$  decays.

Cabibbo  
suppressed

OZI  
suppressed

[PRL. 122\(2019\) 191804; 5.2 fb<sup>-1</sup>](#)



# Observation of $B_{(s)}^0 \rightarrow J/\psi p \bar{p}$ decays

- **First observation of :**

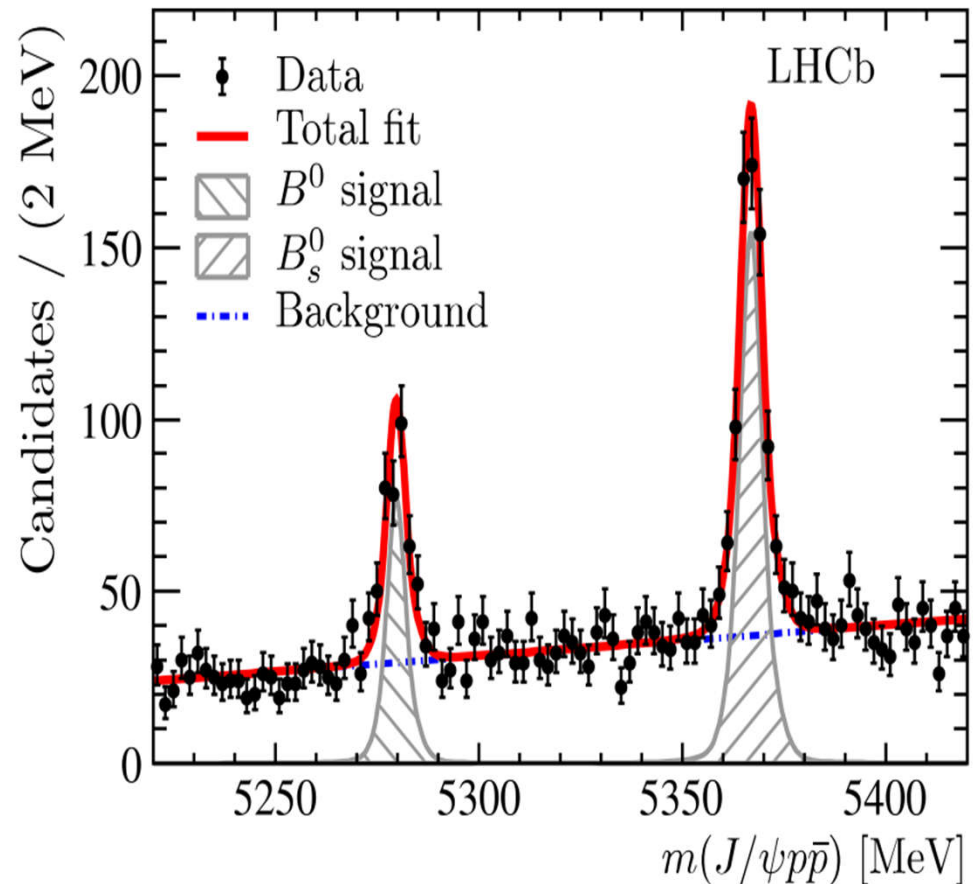
$$B_{(s)}^0 \rightarrow J/\psi p \bar{p}$$

- The measured BRs:

$$\mathcal{B}(B^0 \rightarrow J/\psi p \bar{p}) = (4.51 \pm 0.40 \text{ (stat)} \pm 0.44 \text{ (syst)}) \times 10^{-7},$$

$$\mathcal{B}(B_s^0 \rightarrow J/\psi p \bar{p}) = (3.58 \pm 0.19 \text{ (stat)} \pm 0.39 \text{ (syst)}) \times 10^{-6}.$$

- For the  $B_{(s)}^0$  meson, the result is much higher than the expected value of  $O(10^{-9})$
- Most precise single measurement of both the  $B^0$  and  $B_{(s)}^0$  mass



$$m_{B^0} = 5279.74 \pm 0.30 \text{ (stat)} \pm 0.10 \text{ (syst)} \text{ MeV},$$

$$m_{B_s^0} = 5366.85 \pm 0.19 \text{ (stat)} \pm 0.13 \text{ (syst)} \text{ MeV},$$

**PRL. 122(2019) 191804; 5.2  $fb^{-1}$**



# Observation $\Lambda_b^0 \rightarrow \chi_{c1}(3872)pK^-$

$\Lambda_b^0 \rightarrow \chi_{c1}(3872)pK^-$  with  $\chi_{c1}(3872) \rightarrow J/\psi\pi\pi$ , observed the first time

- Observing  $\Lambda_b^0$  decays involving the  $\chi_{c1}(3872)$  state will allow comparison of their decay rates to the rates for conventional charmonium states
- The BR with respect to that of the  $\Lambda_b^0 \rightarrow \psi(2S)pK^-$ , where the  $\psi(2S) \rightarrow J/\psi\pi\pi$  final state, is measured to be:

$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \chi_{c1}(3872)pK^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \psi(2S)pK^-)} \times \frac{\mathcal{B}(\chi_{c1}(3872) \rightarrow J/\psi\pi^+\pi^-)}{\mathcal{B}(\psi(2S) \rightarrow J/\psi\pi^+\pi^-)} = (5.4 \pm 1.1 \pm 0.2) \times 10^{-2}$$

**JHEP 09 (2019) 028; 1,2,1.9  $fb^{-1}$  at 7,8,13 TeV**

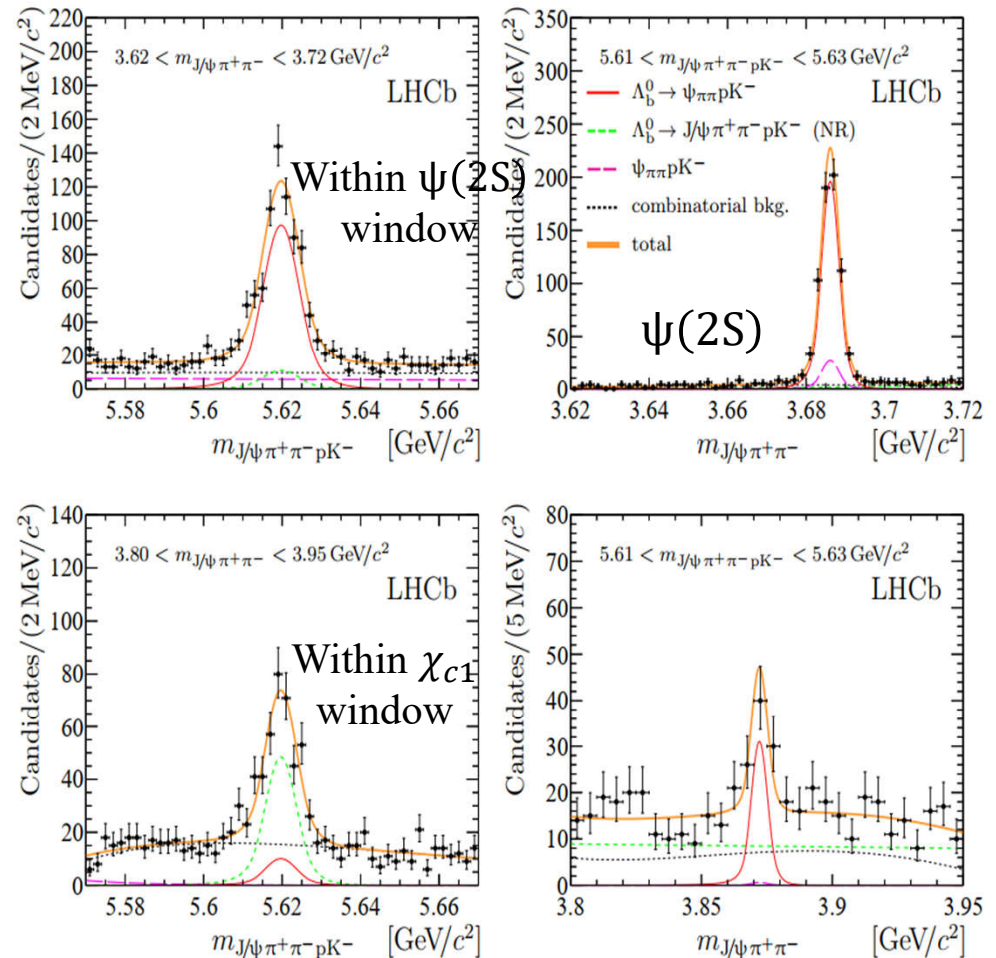
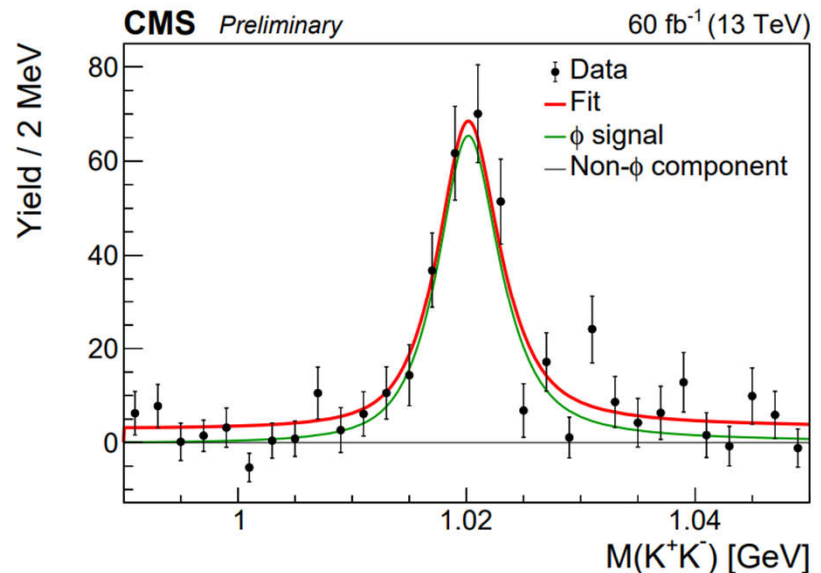
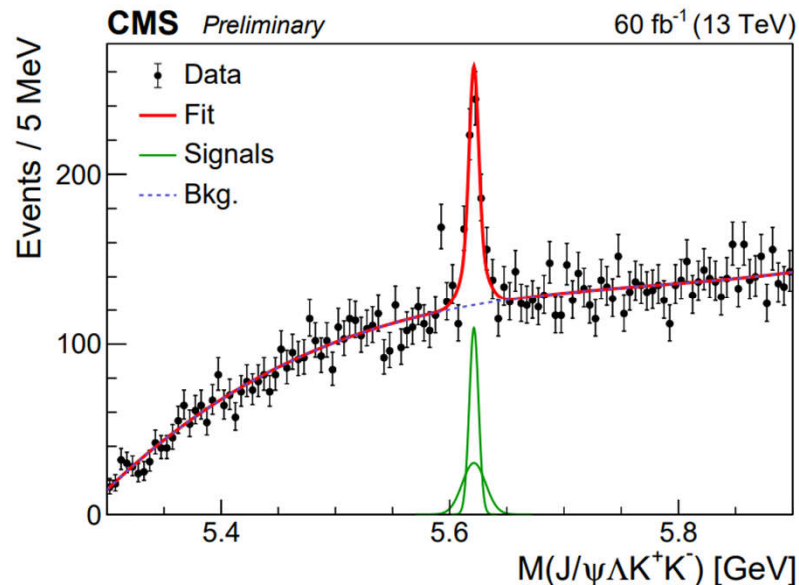


Figure 1: Projection of the two-dimensional distributions of (left)  $J/\psi\pi^+\pi^-pK^-$  and (right)  $J/\psi\pi^+\pi^-$  masses for the (top)  $\Lambda_b^0 \rightarrow \psi(2S)pK^-$  and (bottom)  $\Lambda_b^0 \rightarrow \chi_{c1}(3872)pK^-$  candidates.

# Observation $\Lambda_b \rightarrow J/\psi \Lambda \phi$ (CMS)

$\Lambda_b \rightarrow J/\psi \Lambda \phi$  observed for the first time, at CMS

- Better understanding of final-state strong interactions in the b-baryon decays and test the heavy-quark effective theory
- Rich resonant structure in the  $J/\psi \phi$  system



- **Measured:**

$$\mathcal{B}(\Lambda_b^0 \rightarrow J/\psi \bar{\Lambda} \phi) / \mathcal{B}(\Lambda_b^0 \rightarrow \psi(2S) \Lambda) = (8.26 \pm 0.90 (\text{stat}) \pm 0.68 (\text{syst}) \pm 0.11 (\mathcal{B})) \times 10^{-2}$$

**CMS-PAS-BPH-19-002; 60  $fb^{-1}$  at 13 TeV**

# Observation $B_s^0 \rightarrow \bar{D}^{*0} \phi$ and search $B^0 \rightarrow \bar{D}^0 \phi$

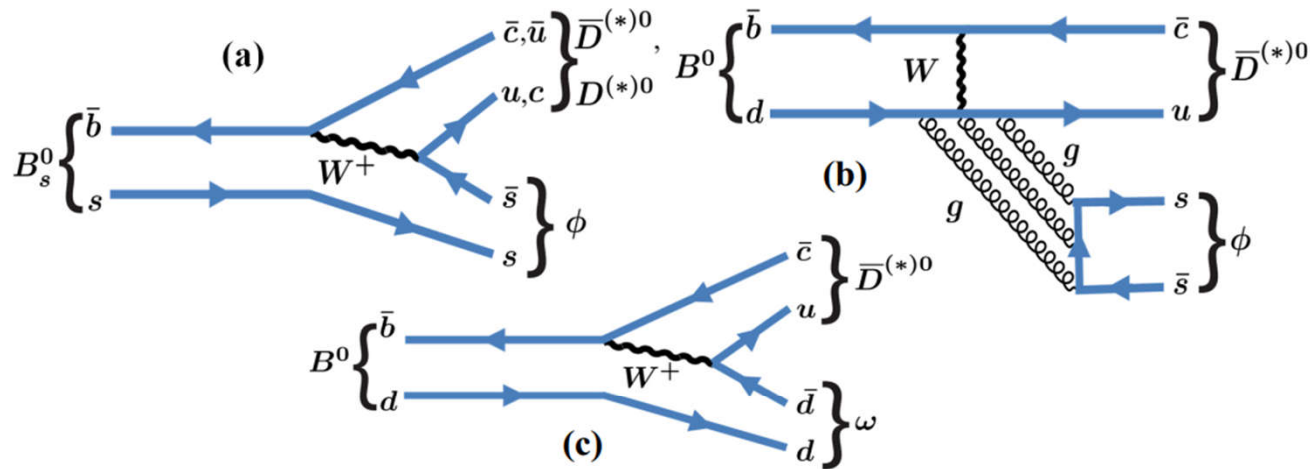
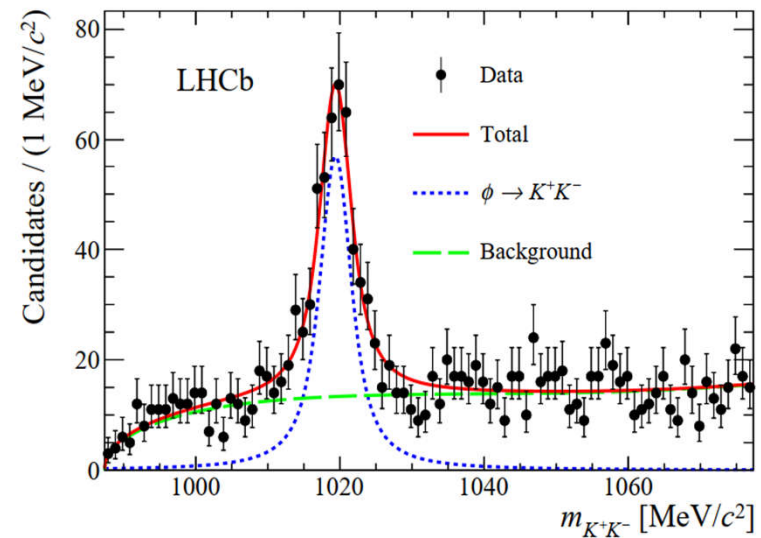


Figure 1: Diagrams that contribute to the (a) colour-suppressed  $B_s^0 \rightarrow \bar{D}^{(*)0}/D^{(*)0} \phi$ , (b)  $W$ -exchange OZI-suppressed  $B^0 \rightarrow \bar{D}^0/D^0 \phi$  and the (c) colour-suppressed  $B^0 \rightarrow \bar{D}^0 \omega$  decays.

- No single measurement dominates the world average of angle  $\gamma$
- Decays  $B_s^0 \rightarrow \bar{D}^{*0} \phi$  open possibilities to offer competitive experimental precision on the angle  $\gamma$



# Observation $B_s^0 \rightarrow \bar{D}^{*0} \phi$ and search $B^0 \rightarrow \bar{D}^0 \phi$

- **First observation of**

$$B_s^0 \rightarrow \bar{D}^{*0} \phi:$$

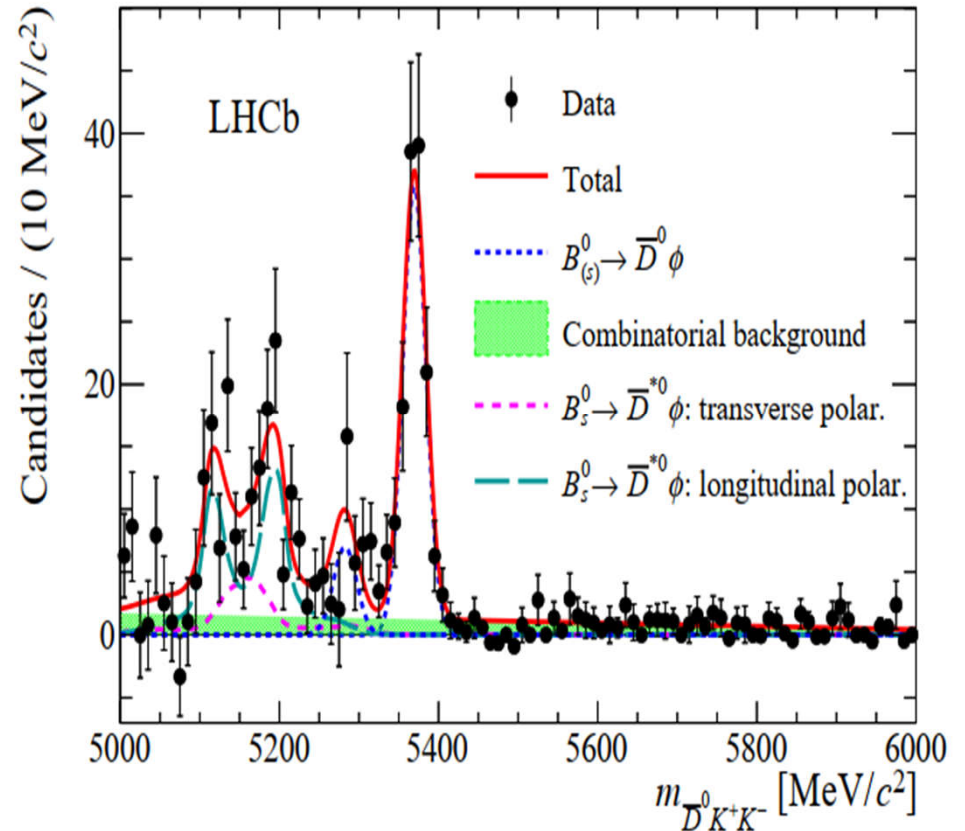
$$\mathcal{B}(B_s^0 \rightarrow \bar{D}^{*0} \phi) = (3.7 \pm 0.5 \pm 0.3 \pm 0.2) \times 10^{-5}$$

- $B^0 \rightarrow \bar{D}^0 \pi^+ \pi^-$  is taken as reference channel

- An upper limit is set:

$$\text{Br}(B^0 \rightarrow \bar{D}^0 \phi) < 2.0(2.3) \times 10^{-6} \text{ at } 90\% (95\%)$$

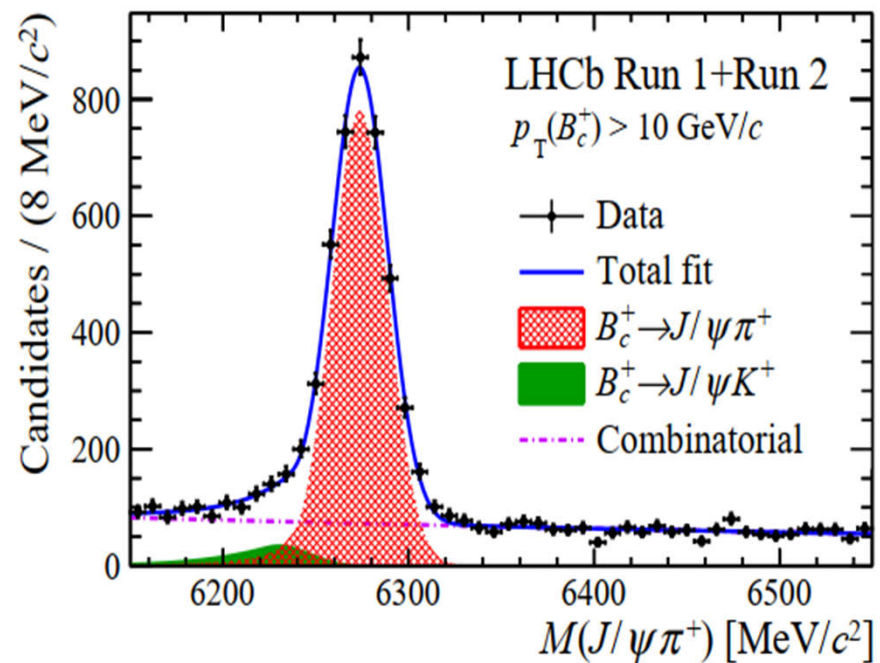
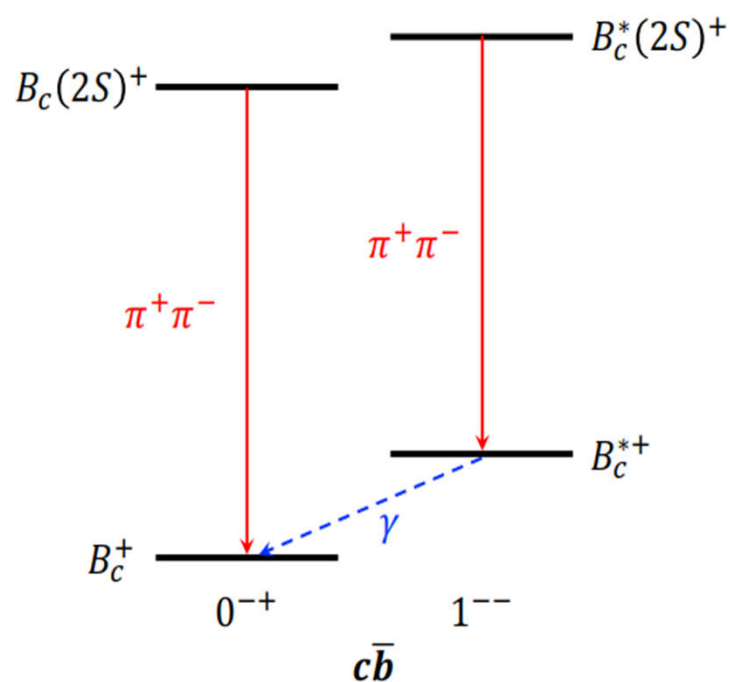
- A factor of six improvement over the previous BaBar result



**PRD (2018) 028**;  $3 \text{ fb}^{-1}$  at 7,8 TeV

# Observation of an excited $B_c^+$ state

- An excited  $B_c^+$  in the  $B_c^+ \pi^+ \pi^-$  invariant-mass:
  - Consistent with the  $B_c^*(2^3S_1)$  state reconstructed without the low-energy  $\gamma$
  - The chain:  $B_c^*(2^3S_1) \rightarrow B_c^*(1^3S_1) \pi^+ \pi^-$ ,  $B_c^*(1^3S_1) \rightarrow B_c^+ \gamma$

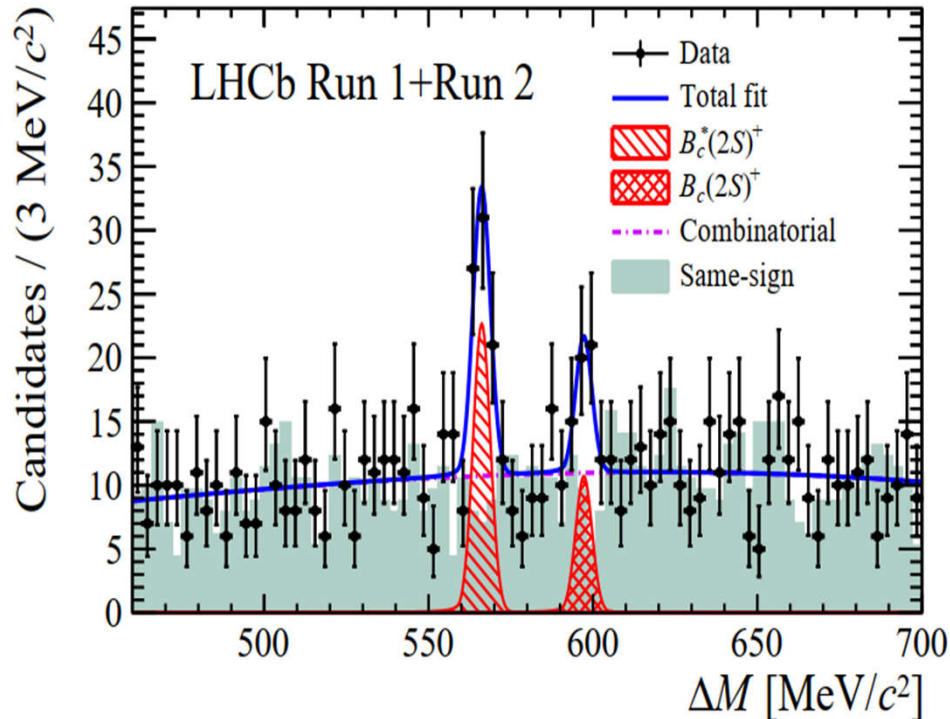


[PRL. 122\(2019\) 232001](#);  $8.5 \text{ fb}^{-1}$  at 7,8,13 TeV

Invariant-mass of the selected  $B_c^+$  candidates  
**Green: K fake as  $\pi$**



# Observation of an excited $B_c^+$ state



Distribution of  $\Delta M = M(B_c^+ \pi^+ \pi^-) - M(B_c^+)$  with the fit results overlaid

	$B_c^*(2S)^+$	$B_c(2S)^+$
Signal yield	$51 \pm 10$	$24 \pm 9$
Peak $\Delta M$ value ( $\text{MeV}/c^2$ )	$566.2 \pm 0.6$	$597.2 \pm 1.3$
Resolution ( $\text{MeV}/c^2$ )	$2.6 \pm 0.5$	$2.5 \pm 1.0$
Local significance	$6.8 \sigma$	$3.2 \sigma$
Global significance	$6.3 \sigma$	$2.2 \sigma$

## First peak:

$6841.2 \pm 0.6$  (stat)  $\pm 0.1$  (syst)  $\pm 0.8$  ( $B_c^+$ )  $\text{MeV}/c^2$ ,

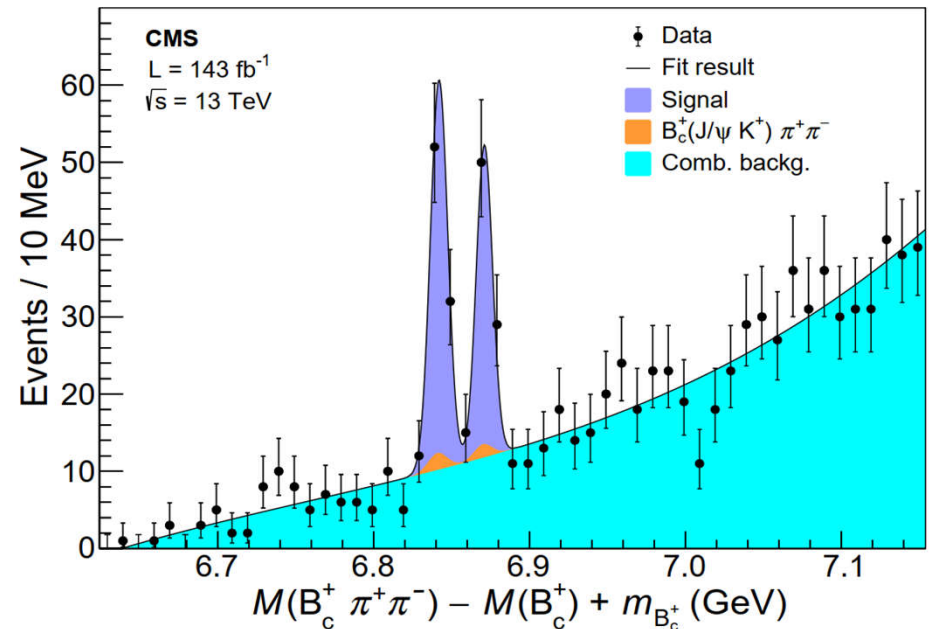
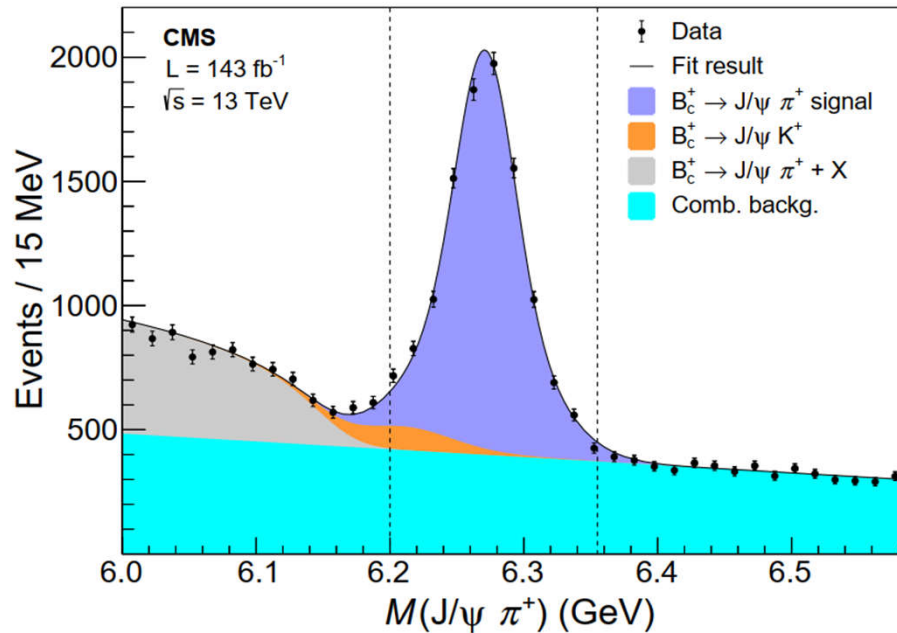
## Second peak:

$6872.1 \pm 1.3$  (stat)  $\pm 0.1$  (syst)  $\pm 0.8$  ( $B_c^+$ )  $\text{MeV}/c^2$ .

[PRL. 122\(2019\) 232001](#);  $8.5 \text{ fb}^{-1}$  at 7,8,13 TeV



# Observation of two excited $B_c^+$ state (CMS)



- **The higher peak mass :**  
 $6871.0 \pm 1.2 \text{ (stat)} \pm 0.8 \text{ (syst)} \pm 0.8 \text{ (} B_c^+ \text{) MeV}$
- **Mass difference: 29MeV**

$M(B_c^+ \pi^+ \pi^-), M(B_c^+)$  :  
 reconstructed  
 $m_{B_c^+}$  : the world-average  $B_c^+$  mass  
**For better resolution**

**PRL. 122(2019) 132001; 13 fb<sup>-1</sup> at 13 TeV**

# $B^0 \rightarrow \eta_c(1S)K^+\pi^-$

- Theory models predict  $0^+$  candidates below the open-charm threshold that could decay into  $\eta_c\pi^-$  (e.g. the  $Z_c(3900)^-$ ).
- Therefore discovery of a charged charmonium-like meson in  $(\eta_c\pi^-) \Rightarrow$  important input towards understanding the nature of exotic hadrons.
- $B^0 \rightarrow \eta_c(1S)K^+\pi^-$  decay is studied for the first time

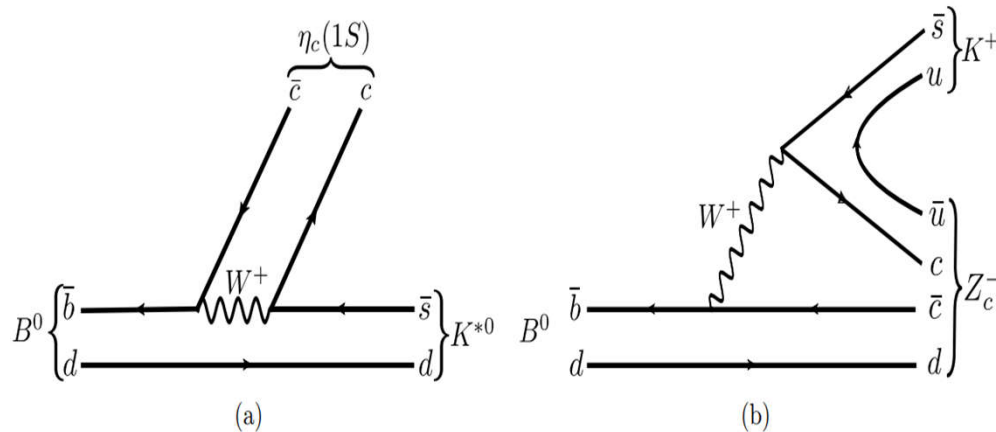
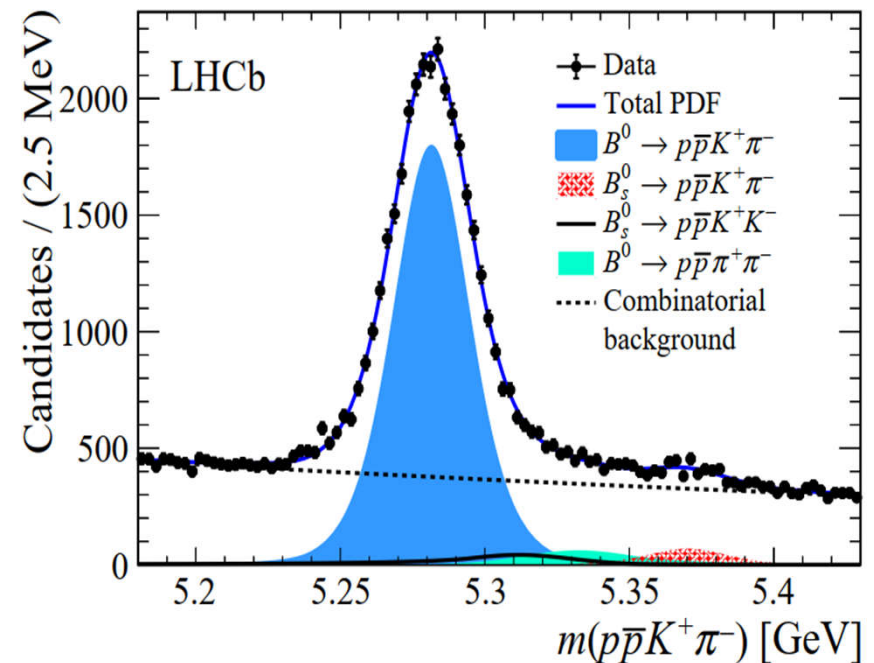


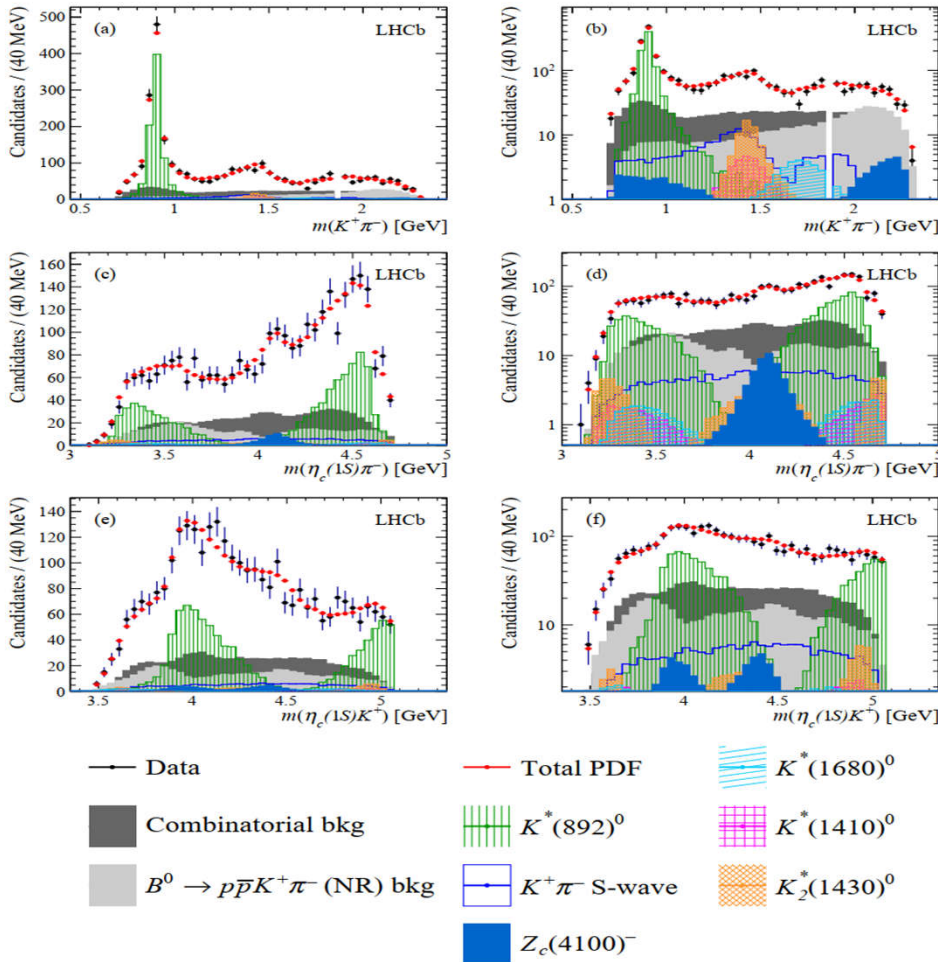
Figure 1: Feynman diagrams for (a)  $B^0 \rightarrow \eta_c K^{*0}$  and (b)  $B^0 \rightarrow Z_c^- K^+$  decay sequences.

**EPJC78 (2018) 1019**;  $4.7 fb^{-1}$  at 7,8,13 TeV



# $B^0 \rightarrow \eta_c(1S)K^+\pi^-$

- Good description is obtained when including an exotic  $Z_c(\eta_c(1S)\pi^-)$  resonant state ( $> 3\sigma$ )



- First measurement of  $B^0 \rightarrow \eta_c(1S)K^+\pi^-$  BR:

$$\mathcal{B}(B^0 \rightarrow \eta_c K^+ \pi^-) = (5.73 \pm 0.24 \pm 0.13 \pm 0.66) \times 10^{-4},$$

- For the resonance:

$$m_{Z_c^-} = 4096 \pm 20^{+18}_{-22} \text{ MeV and } \Gamma_{Z_c^-} = 152 \pm 58^{+60}_{-35} \text{ MeV}$$

- **The fitted fractions:**

Amplitude	Fit fraction (%)
$B^0 \rightarrow \eta_c K^*(892)^0$	$51.4 \pm 1.9^{+1.7}_{-4.8}$
$B^0 \rightarrow \eta_c K^*(1410)^0$	$2.1 \pm 1.1^{+1.1}_{-1.1}$
$B^0 \rightarrow \eta_c K^+ \pi^-$ (NR)	$10.3 \pm 1.4^{+1.0}_{-1.2}$
$B^0 \rightarrow \eta_c K_0^*(1430)^0$	$25.3 \pm 3.5^{+3.5}_{-2.8}$
$B^0 \rightarrow \eta_c K_2^*(1430)^0$	$4.1 \pm 1.5^{+1.0}_{-1.6}$
$B^0 \rightarrow \eta_c K^*(1680)^0$	$2.2 \pm 2.0^{+1.5}_{-1.7}$
$B^0 \rightarrow \eta_c K_0^*(1950)^0$	$3.8 \pm 1.8^{+1.4}_{-2.5}$
$B^0 \rightarrow Z_c(4100)^- K^+$	$3.3 \pm 1.1^{+1.2}_{-1.1}$

# Measurements of $\Xi_b^-$ baryons

- **The first measurement** of the production rate of  $\Xi_b^-$  baryons in  $pp$  collisions
  - Relative to that of  $\Lambda_b^0$  baryons

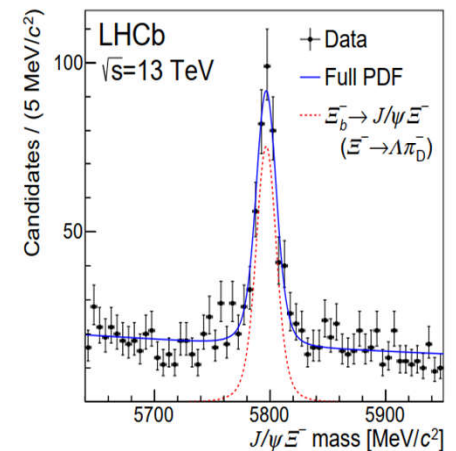
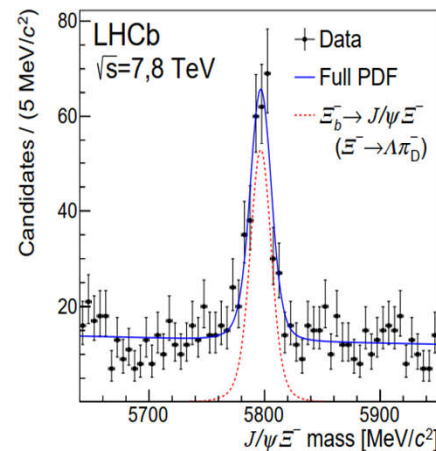
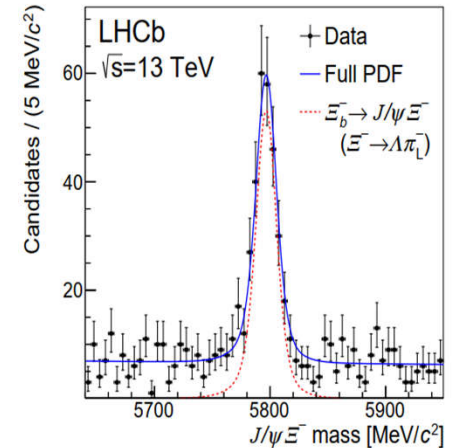
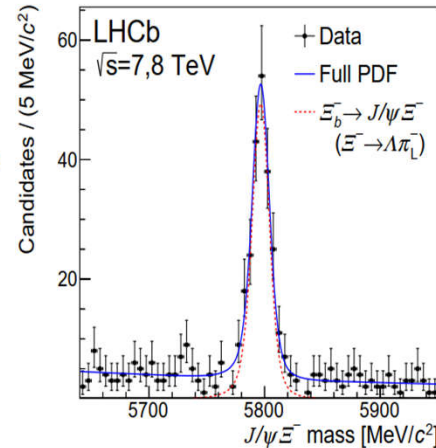
$$\frac{f_{\Xi_b^-} \mathcal{B}(\Xi_b^- \rightarrow J/\psi \Xi^-)}{f_{\Lambda_b^0} \mathcal{B}(\Lambda_b^0 \rightarrow J/\psi \Lambda)} = (10.8 \pm 0.9 \pm 0.8) \times 10^{-2} \quad [\sqrt{s} = 7,8 \text{ TeV}],$$

$$\frac{f_{\Xi_b^-} \mathcal{B}(\Xi_b^- \rightarrow J/\psi \Xi^-)}{f_{\Lambda_b^0} \mathcal{B}(\Lambda_b^0 \rightarrow J/\psi \Lambda)} = (13.1 \pm 1.1 \pm 1.0) \times 10^{-2} \quad [\sqrt{s} = 13 \text{ TeV}],$$

- The mass of the  $\Xi_b^-$  is also measured relative to that of the  $\Lambda_b^0$  baryon:

$$m(\Xi_b^-) = 5796.70 \pm 0.39 \pm 0.15 \pm 0.17 \text{ MeV}/c^2,$$

- The last uncertainty is due to the precision on the known  $\Lambda_b^0$  mass
- Most precise determination of  $\Xi_b^-$  mass.



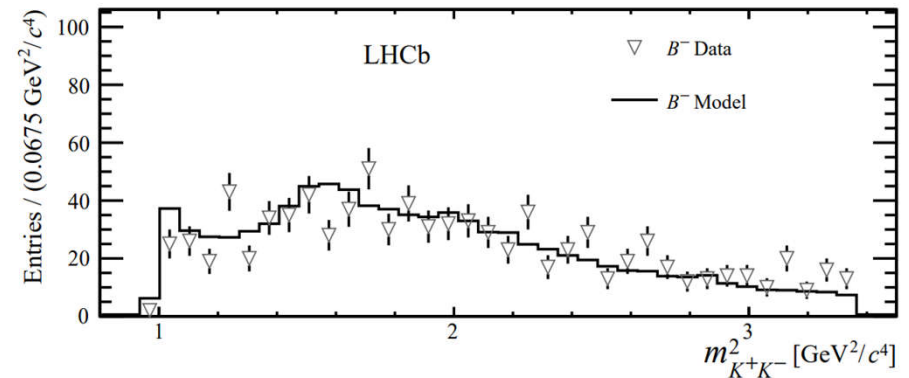
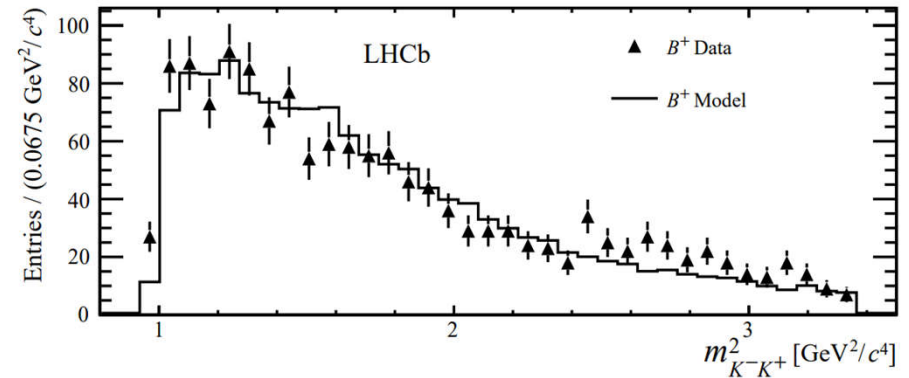
**PRD 99052006 (2019); 1,2,1.6 fb<sup>-1</sup> at 7,8,13 TeV**



# Amplitude analysis of $B^{\pm} \rightarrow \pi^{\pm} K^+ K^-$

- The first amplitude analysis of the  $B^{\pm} \rightarrow \pi^{\pm} K^+ K^-$
- The data is best described by a coherent sum of **five resonant structures + a nonresonant component + a  $\pi\pi \leftrightarrow KK$  S-wave rescattering**

Contribution	Fit Fraction(%)	$A_{CP}(\%)$	Magnitude ( $B^+/B^-$ )	Phase $^{\circ}$ ( $B^+/B^-$ )
$K^*(892)^0$	$7.5 \pm 0.6 \pm 0.5$	$+12.3 \pm 8.7 \pm 4.5$	$0.94 \pm 0.04 \pm 0.02$	0 (fixed)
			$1.06 \pm 0.04 \pm 0.02$	0 (fixed)
$K_0^*(1430)^0$	$4.5 \pm 0.7 \pm 1.2$	$+10.4 \pm 14.9 \pm 8.8$	$0.74 \pm 0.09 \pm 0.09$	$-176 \pm 10 \pm 16$
			$0.82 \pm 0.09 \pm 0.10$	$136 \pm 11 \pm 21$
Single pole	$32.3 \pm 1.5 \pm 4.1$	$-10.7 \pm 5.3 \pm 3.5$	$2.19 \pm 0.13 \pm 0.17$	$-138 \pm 7 \pm 5$
			$1.97 \pm 0.12 \pm 0.20$	$166 \pm 6 \pm 5$
$\rho(1450)^0$	$30.7 \pm 1.2 \pm 0.9$	$-10.9 \pm 4.4 \pm 2.4$	$2.14 \pm 0.11 \pm 0.07$	$-175 \pm 10 \pm 15$
			$1.92 \pm 0.10 \pm 0.07$	$140 \pm 13 \pm 20$
$f_2(1270)$	$7.5 \pm 0.8 \pm 0.7$	$+26.7 \pm 10.2 \pm 4.8$	$0.86 \pm 0.09 \pm 0.07$	$-106 \pm 11 \pm 10$
			$1.13 \pm 0.08 \pm 0.05$	$-128 \pm 11 \pm 14$
Rescattering	$16.4 \pm 0.8 \pm 1.0$	$-66.4 \pm 3.8 \pm 1.9$	$1.91 \pm 0.09 \pm 0.06$	$-56 \pm 12 \pm 18$
			$0.86 \pm 0.07 \pm 0.04$	$-81 \pm 14 \pm 15$
$\phi(1020)$	$0.3 \pm 0.1 \pm 0.1$	$+9.8 \pm 43.6 \pm 26.6$	$0.20 \pm 0.07 \pm 0.02$	$-52 \pm 23 \pm 32$
			$0.22 \pm 0.06 \pm 0.04$	$107 \pm 33 \pm 41$

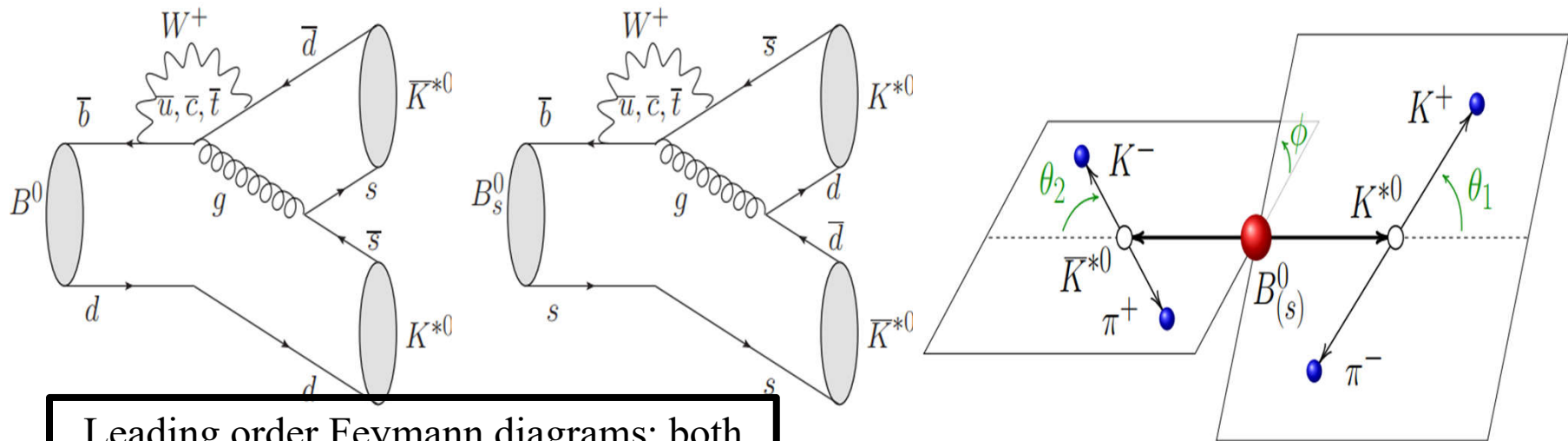


← The fitted fractions

**ARXIV:1905.09244(to PRL);  $3fb^{-1}$**

# Amplitude analysis of $B_s^0 \rightarrow K^{*0} K^{*0}$

- $B_s^0 \rightarrow K^{*0} \bar{K}^{*0}$  could be a golden channel for a precision test of the CKM  $\beta_s$ .
- High-precision analyses of it, require to account for subleading amplitudes



Leading order Feynman diagrams; both are dominated by gluonic-penguin

## The amplitude analysis outputs:

- Measured observables are compatible with the absence of  $CP$  violation
- A low polarisation fraction is found
- A large S-wave contribution (60%), is measured in the 150 MeV window around the  $K^{*0}$  mass.

[JHEP 07 \(2019\) 032](#);  $3 \text{ fb}^{-1}$



# Measurement of $B^0 \rightarrow K^{*0} \bar{K}^{*0}$

- The first study of  $B^0 \rightarrow K^{*0} \bar{K}^{*0}$  decays is performed

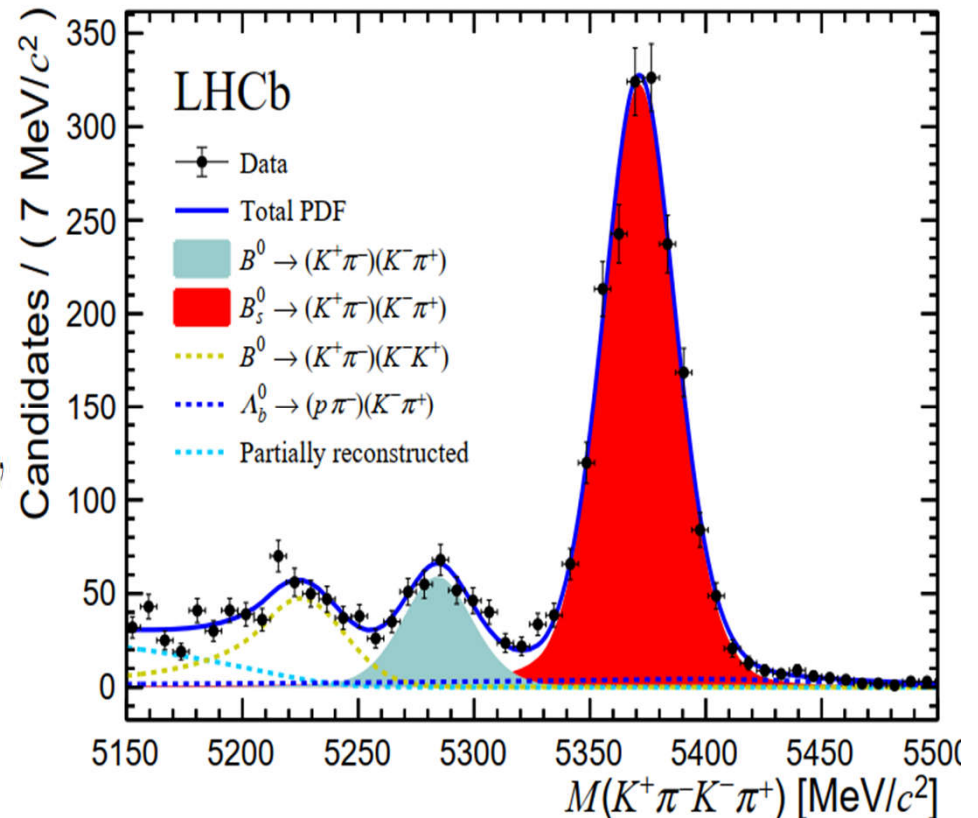
- The ratio of BR:

$$\frac{\mathcal{B}(B^0 \rightarrow K^{*0} \bar{K}^{*0})}{\mathcal{B}(B_s^0 \rightarrow K^{*0} \bar{K}^{*0})} = 0.0758 \pm 0.0057 (\text{stat}) \pm 0.0025 (\text{syst}) \pm 0.0016 \left( \frac{f_s}{f_d} \right),$$

- With this the  $B^0 \rightarrow K^{*0} \bar{K}^{*0}$  BR is found to be:

$$\mathcal{B}(B^0 \rightarrow K^{*0} \bar{K}^{*0}) = (8.0 \pm 0.9 (\text{stat}) \pm 0.4 (\text{syst})) \times 10^{-7}.$$

**JHEP 07 (2019) 032;  $3 \text{ fb}^{-1}$**



Aggregated four-body invariant-mass fit result of the 2011 and 2012 data.

# $B_s^0 \rightarrow K_s^0 K^\pm \pi^\mp$ amplitude analysis

- Undiscovered particles could contribute in the  $b \rightarrow s$  loops and cause the observables to deviate from the values expected in the SM

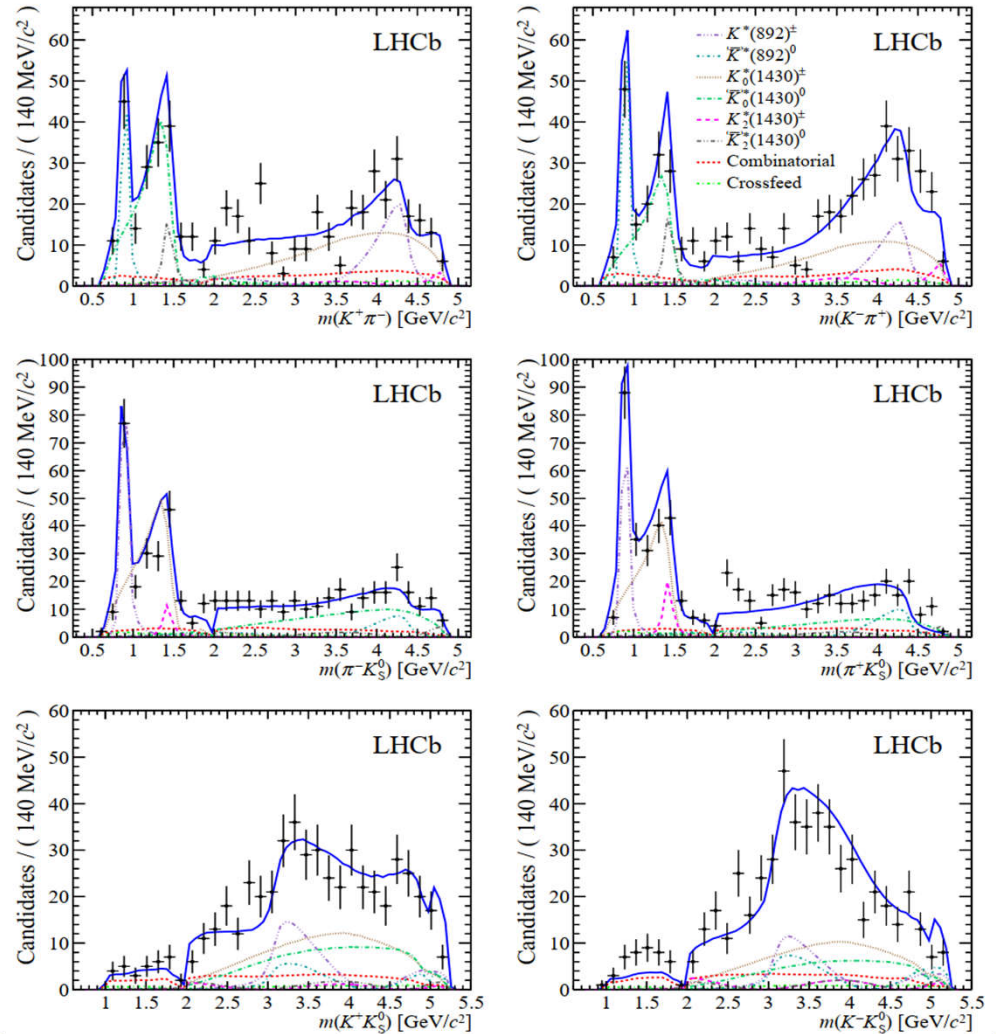
$K_s^0 K^+ \pi^-$		$K_s^0 K^- \pi^+$	
Resonance	Fit fraction (%)	Resonance	Fit fraction (%)
$K^*(892)^-$	$15.6 \pm 1.5$	$K^*(892)^+$	$13.4 \pm 2.0$
$K_0^*(1430)^-$	$30.2 \pm 2.6$	$K_0^*(1430)^+$	$28.5 \pm 3.6$
$K_2^*(1430)^-$	$2.9 \pm 1.3$	$K_2^*(1430)^+$	$5.8 \pm 1.9$
$K^*(892)^0$	$13.2 \pm 2.4$	$\bar{K}^*(892)^0$	$19.2 \pm 2.3$
$K_0^*(1430)^0$	$33.9 \pm 2.9$	$\bar{K}_0^*(1430)^0$	$27.0 \pm 4.1$
$K_2^*(1430)^0$	$5.9 \pm 4.0$	$\bar{K}_2^*(1430)^0$	$7.7 \pm 2.8$

The fit fractions associated with each resonant component

- The decays are observed for the first time:

$$B_s^0 \rightarrow K_0^*(1430)^\pm K^\mp \text{ and } B_s^0 \rightarrow \bar{K}_0^*(1430)^0 \bar{K}^0$$

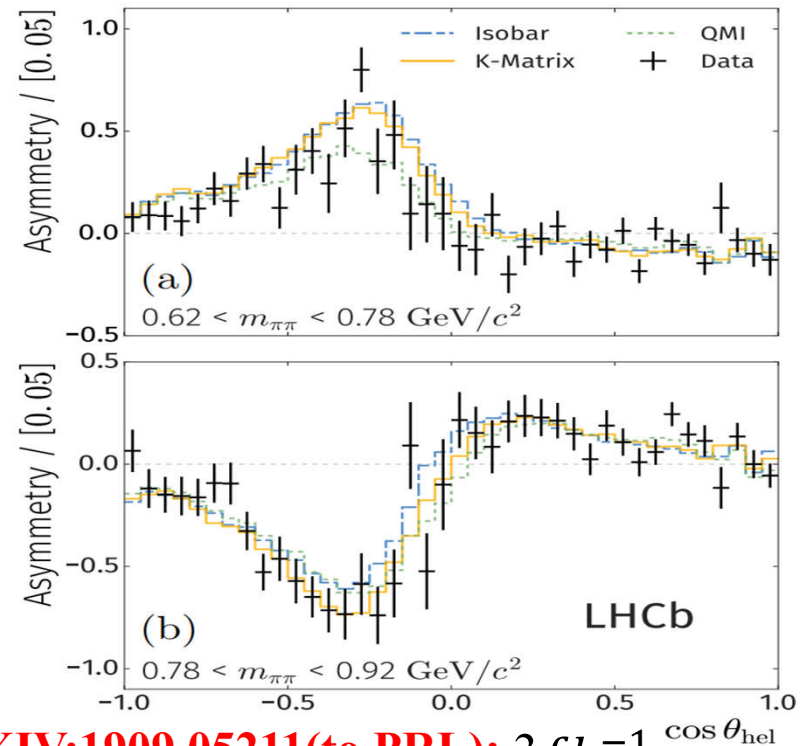
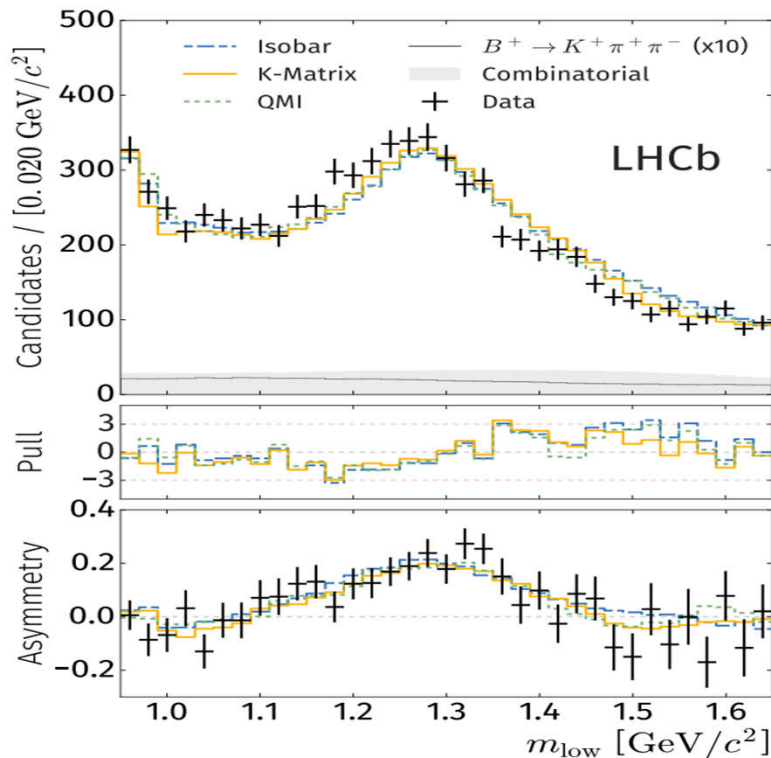
**JHEP 06(2019) 114;  $3 \text{ fb}^{-1}$**



# Several sources of CPV in $B^+ \rightarrow \pi^+ \pi^+ \pi^-$

**Observations of CPV** from an amplitude analysis:

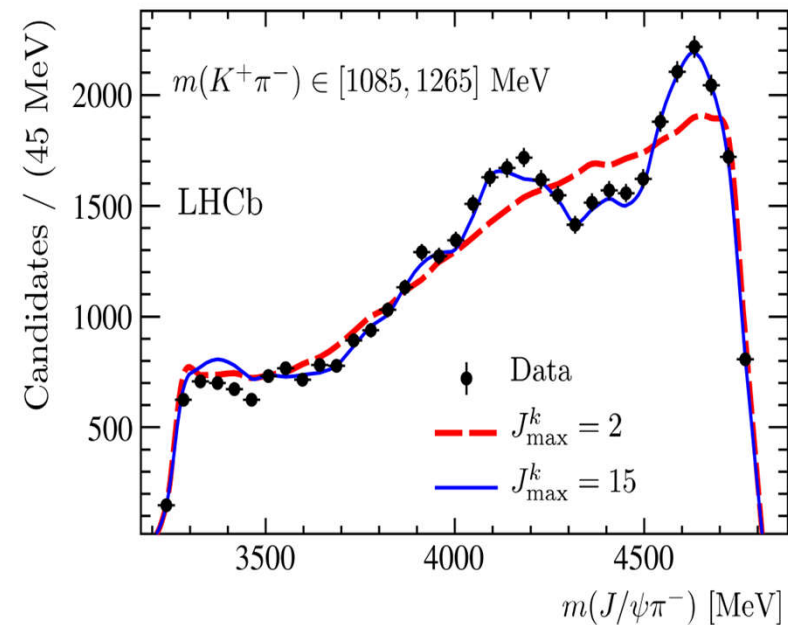
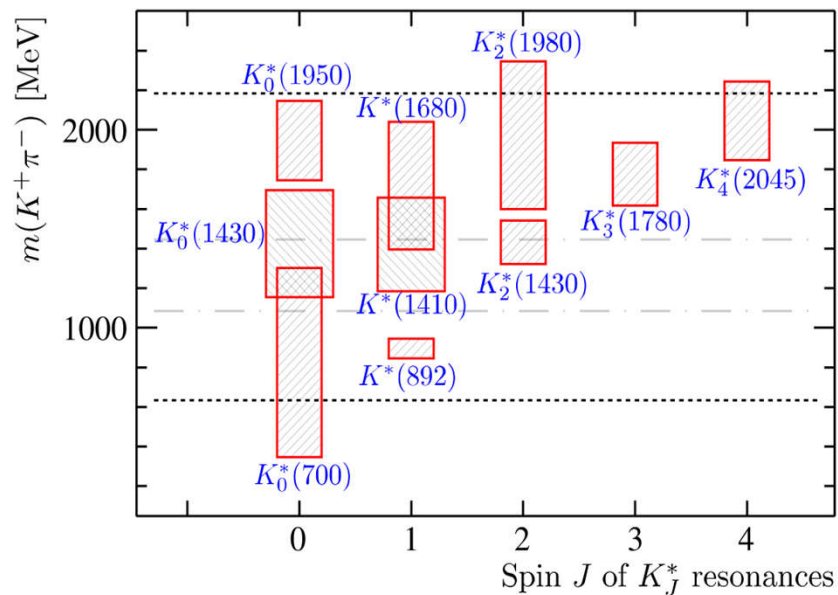
- A large CP asymmetry is observed in the decay amplitude involving the tensor  $f_2(1270)$  resonance
- Significant CPV is found in the  $\pi^+ \pi^-$  S-wave at low invariant mass.
- CPV related to interference between the  $\pi^+ \pi^-$  S-wave and the P-wave  $B^+ \rightarrow \rho(770)^0 \pi^+$  amplitude is also established



[ARXIV:1909.05211\(to PRL\)](https://arxiv.org/abs/1909.05211);  $3fb^{-1}$

# Model-ind. observation of $B^0 \rightarrow J/\psi K^+ \pi^-$

- $K^+ \pi^-$  spectrum into 5 bins. In each bin, the hypothesis that the 3-D angular distribution can be described by structures induced only by  $K^*$  resonances is examined, making minimal assumptions about the  $K^+ \pi^-$  system.
- Data reject the  $K^*$ -only hypothesis, implying the observation of exotic contributions in a model independent fashion.
- $m(J/\psi \pi^-)$  vs.  $m(K^+ \pi^-)$  suggests structures near  $m(J/\psi \pi^-) = 4200 \text{ MeV}, 4600 \text{ MeV}$



**PRL. 122(2019) 152002;  $3 \text{ fb}^{-1}$**

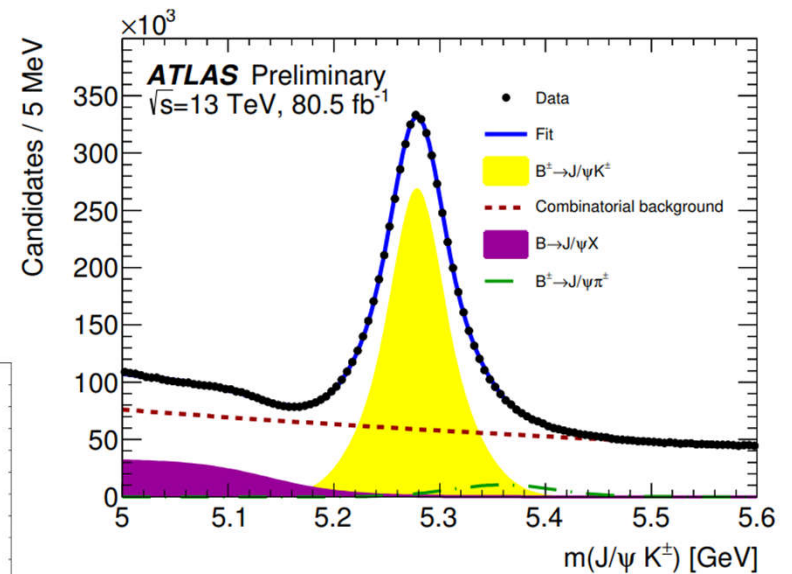
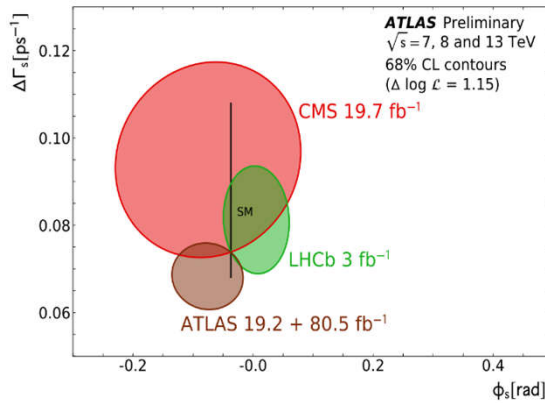
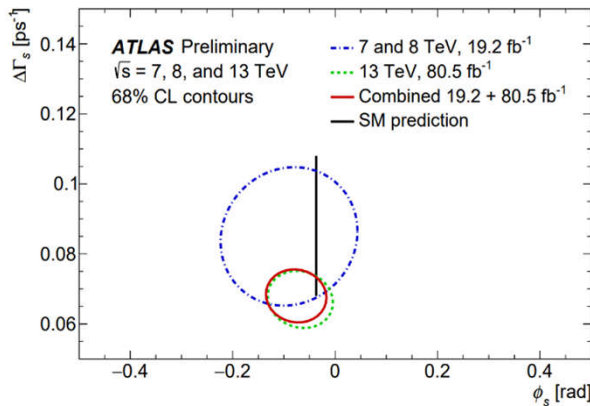
$J_{\max k} = 2$  model clearly misses the peaking structures in the data around 4200, 4600 MeV



# CPV phase $\phi_s$ in $B_S^0 \rightarrow J/\psi\phi$ (ATLAS)

- The CP violating phase  $\phi_s$  is defined as the weak phase difference between the  $B_S^0 - \bar{B}_S^0$  mixing amplitude and the  $b \rightarrow ccs$  decay amplitude
- NP involved in the mixing may increase by enhancing the mixing phase  $\phi_s$  with respect to the SM
- Combined 7,8 TeV data, **gives the most stringent measurements**

$$\begin{aligned} \phi_s &= -0.076 \pm 0.034 \text{ (stat.)} \pm 0.019 \text{ (syst.) rad} \\ \Delta\Gamma_s &= 0.068 \pm 0.004 \text{ (stat.)} \pm 0.003 \text{ (syst.) ps}^{-1} \\ \Gamma_s &= 0.669 \pm 0.001 \text{ (stat.)} \pm 0.001 \text{ (syst.) ps}^{-1} \end{aligned}$$



**Flavour tagging of B meson:**  
 Calibrations from  
 $B^\pm \rightarrow J/\psi K^\pm$ , applied to  
 $B_S^0 \rightarrow J/\psi(\mu^+\mu^-)\phi(K^+K^-)$

**ATLAS-CONF-2019-009; 80.5 fb<sup>-1</sup> at 13TeV**

# Summary

- Hadronic B decays play a key role in checking the SM. Numerous new results are out since PIC 2018, only part of them are shown today.
- We have much more Run2 data to analyze, more results are coming
- LHCb currently ongoing a major upgrade for Run-3 and Run-4
  - Preparations underway for a new era of discoveries taking maximum advantage of the High-Luminosity LHC



# **Backup Slides**

# Overview of the timeline

- LHC Run-I (2010-2013) & LHC Run-II (2015-2018)
- LHC Run-III, Run-IV (2021-2023, 2026-2029)
  - based on LHCb Upgrade [I(a), I(b)]
- LHC Run-V (2031-)
  - based on LHCb Upgrade II

